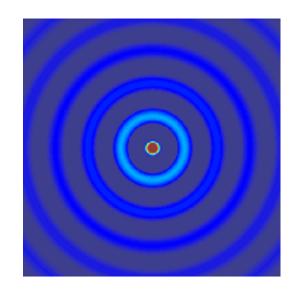
ECE 6341

Spring 2016

Prof. David R. Jackson ECE Dept.



Notes 33

Steepest-Descent Method

Complex Integral:

$$I(\Omega) = \int_{C} f(z)e^{\Omega g(z)}dz$$

This is an extension of Laplace's method to treat integrals in the *complex plane*.

The method was published by Peter Debye in 1909. Debye noted in his work that the method was developed in a unpublished note by Bernhard Riemann (1863).



Peter Joseph William Debye (March 24, 1884 – November 2, 1966) was a Dutch physicist and physical chemist, and Nobel laureate in Chemistry.



Georg Friedrich Bernhard
Riemann (September 17, 1826 –
July 20, 1866) was an influential
German mathematician who made
lasting contributions to analysis
and differential geometry, some of
them enabling the later
development of general relativity.

http://en.wikipedia.org/wiki/Peter_Debye

http://en.wikipedia.org/wiki/Bernhard_Riemann

Complex Integral:

$$I(\Omega) = \int_{C} f(z) e^{\Omega g(z)} dz$$

The functions f(z) and g(z) are analytic (except for poles or branch points), so that the path C may be deformed if necessary (possibly adding residue contributions or branch-cut integrals).

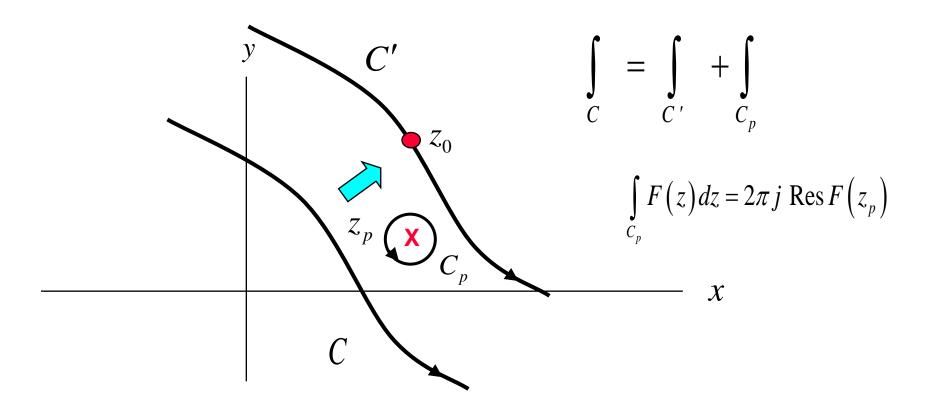
Saddle Point (SP):
$$g'(z_0) = 0$$

$$\frac{\partial g}{\partial x} = 0 \qquad \frac{\partial g}{\partial y} = 0$$

Path deformation:

If the path does not go through a saddle point, we assume that it can be deformed to do so.

If any singularities are encountered during the path deformation, they must be accounted for (e.g., residue of captured poles).



Denote
$$g(z) = u(z) + jv(z)$$

Cauchy Reimann eqs.:
$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}$$
$$\frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}$$

Hence
$$\frac{\partial^2 u}{\partial x^2} = \frac{\partial}{\partial x} \left(\frac{\partial u}{\partial x} \right) = \frac{\partial}{\partial x} \left(\frac{\partial v}{\partial y} \right) = \frac{\partial}{\partial y} \left(\frac{\partial v}{\partial x} \right)$$
$$= -\frac{\partial^2 u}{\partial y^2}$$

or
$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

$$\longrightarrow \quad \text{If} \quad u_{xx} < 0, \quad \text{then} \quad u_{yy} > 0$$

Near the saddle point:

$$u(x, y) \approx u(x_0, y_0) + \frac{1}{2}u_{xx}(x - x_0)^2 + \frac{1}{2}u_{yy}(y - y_0)^2 + u_{xy}(x - x_0)(y - y_0)$$

Ignore (rotate coordinates to eliminate).

$$u(x,y) \approx u(x_0,y_0) + \frac{1}{2}u_{xx}(x-x_0)^2 + \frac{1}{2}u_{yy}(y-y_0)^2 + u_{xy}(x-x_0)(y-y_0)$$

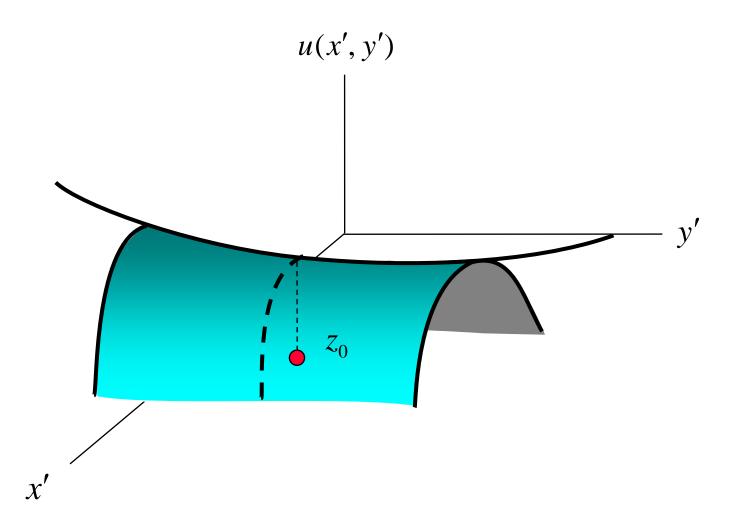
In the rotated coordinate system:

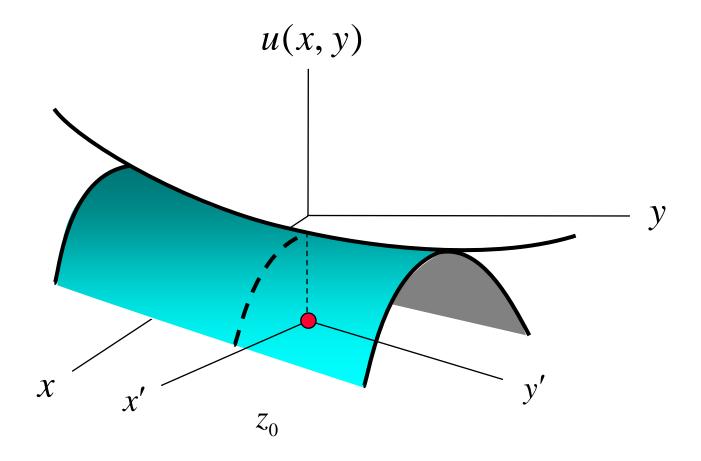
$$u(x', y') \approx u(x'_0, y'_0) + \frac{1}{2}u_{x'x'}(x' - x'_0)^2 + \frac{1}{2}u_{y'y'}(y' - y'_0)^2$$

Assume that the coordinate system is rotated so that

$$u_{x'x'} < 0 \qquad u_{y'y'} > 0$$

The u(x', y') function has a "saddle" shape near the saddle point:





Note: The saddle does not necessarily open along one of the principal axes (only when $u_{xy}(x_0, y_0) = 0$).

Along any <u>descending path</u> (where the *u* function decreases):

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_C e^{\Omega[g(z)-g(z_0)]} dz$$

or

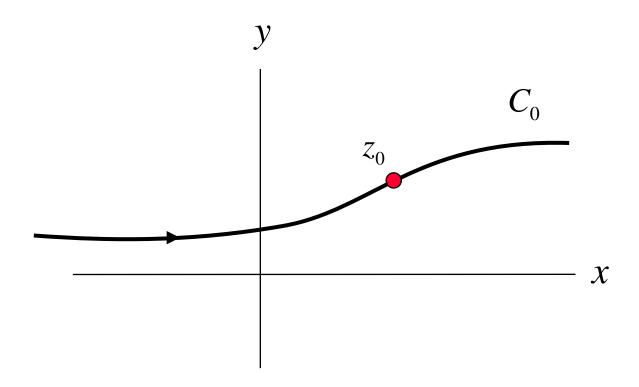
$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_C e^{j\Omega[v(z)-v(z_0)]} e^{\Omega[u(z)-u(z_0)]} dz$$
Behaves like a delta function

Both the phase and amplitude change along an arbitrary descending path *C*.

If we can find a path along which the phase does not change (ν is constant), then the integral will look like that in Laplace's method.

Choose path of constant phase:

$$C_0: v(z) = v(z_0) = \text{constant}$$



Gradient Property (proof on next slide):

$$\nabla u(x,y)$$
 is parallel to C_0

Hence C_0 is either a "path of steepest descent" (SDP) or a "path of steepest ascent" (SAP).

(Of course, we want to choose the SDP.)

SDP: u(x,y) decreases as fast as possible along the path away from the saddle point.

SAP: u(x,y) increases as fast as possible along the path away from the saddle point.

Proof of gradient property

$$g(z) = u(z) + jv(z)$$

$$\nabla u = \hat{\underline{x}} \frac{\partial u}{\partial x} + \hat{\underline{y}} \frac{\partial u}{\partial y}$$

$$\nabla v = \hat{\underline{x}} \frac{\partial v}{\partial x} + \hat{\underline{y}} \frac{\partial v}{\partial y}$$

$$\nabla u \cdot \nabla v = \frac{\partial u}{\partial x} \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \frac{\partial v}{\partial y}$$
$$= \frac{\partial u}{\partial x} \left(-\frac{\partial u}{\partial y} \right) + \frac{\partial u}{\partial y} \left(\frac{\partial u}{\partial x} \right)$$
$$= 0$$

Hence, $\nabla u \perp \nabla v$

Also, $\nabla v \perp C_0$ (v is constant on C_0)

 $\begin{array}{c|c}
\hline
C_0 \\
\hline
\nabla u \\
\hline
\end{array}$

Hence
$$\nabla u \parallel C_0$$

Recall:

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_C e^{j\Omega[v(z)-v(z_0)]} e^{\Omega[u(z)-u(z_0)]} dz$$

Because the *v* function is constant along the SDP, we have

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_{SDP} e^{\Omega[u(z)-u(z_0)]} dz$$

or

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_{SDP} e^{\Omega \left[g(z) - g(z_0)\right]} dz$$

This is real on the SDP.

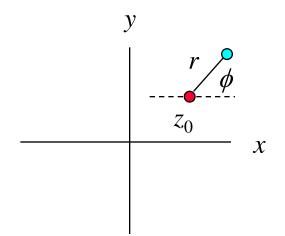
Local behavior near the saddle point

$$g(z) \approx g(z_0) + g'(z_0)(z-z_0) + \frac{1}{2}g''(z_0)(z-z_0)^2$$

so
$$g(z) - g(z_0) \approx \frac{1}{2}g''(z_0)(z - z_0)^2$$

Denote
$$g''(z_0) = R e^{j\alpha}$$

$$z - z_0 = re^{j\phi}$$



Then we have

$$g(z)-g(z_0) \approx \frac{1}{2} R r^2 e^{j(\alpha+2\phi)}$$

$$g(z)-g(z_0) \approx \frac{1}{2}(Rr^2)e^{j(\alpha+2\phi)} \quad \Longrightarrow \quad u(z)-u(z_0) \approx \frac{1}{2}(Rr^2)\cos(\alpha+2\phi)$$

SAP:
$$\alpha + 2\phi = 0 + 2\pi n$$

$$\phi = -\frac{\alpha}{2}, -\frac{\alpha}{2} + \pi$$

SDP:
$$\alpha + 2\phi = \pi + 2\pi n$$

$$\phi = -\frac{\alpha}{2} - \frac{\pi}{2}, -\frac{\alpha}{2} - \frac{\pi}{2} + \pi$$

$$u(z) - u(z_0) = \frac{1}{2}Rr^2$$
$$v(z) - v(z_0) = 0$$

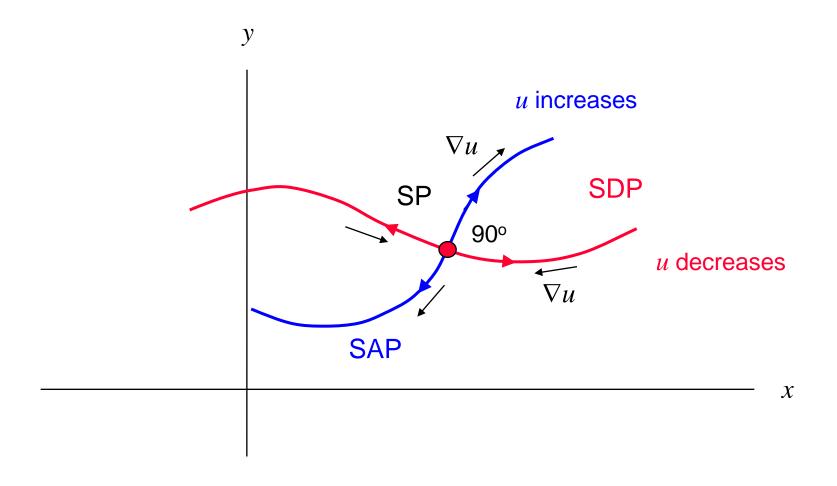
$$\alpha + 2\phi = \pi + 2\pi n$$

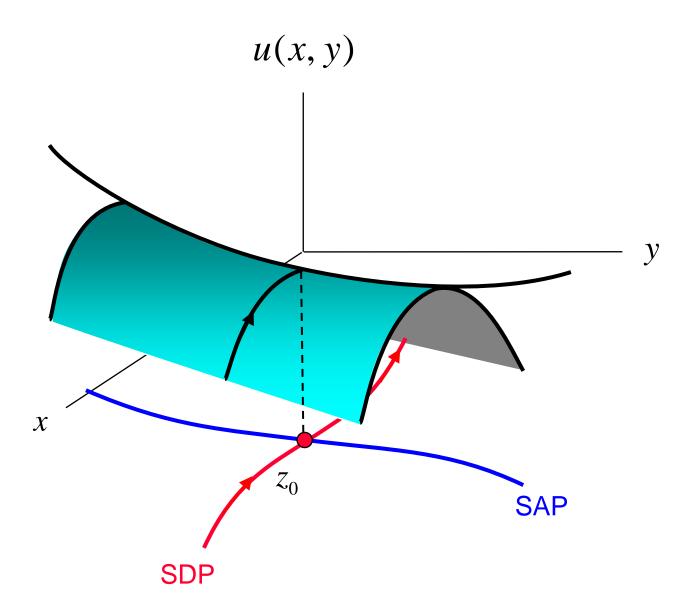
$$\phi = -\frac{\alpha}{2} - \frac{\pi}{2}, -\frac{\alpha}{2} - \frac{\pi}{2} + \pi$$

$$u(z) - u(z_0) = -\frac{1}{2}Rr^2$$

$$v(z) - v(z_0) = 0$$

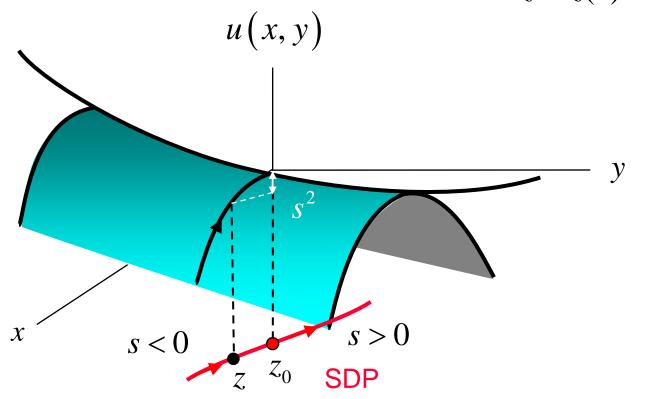
Note: The two paths are 90° apart at the saddle point.





$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_{SDP} e^{\Omega[u(z)-u(z_0)]} dz$$

Define
$$s^2 \equiv u(z_0) - u(z)$$
 This defines $z = z(s)$



$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \int_{-\infty}^{+\infty} e^{-\Omega s^2} \left(\frac{dz}{ds}\right) ds$$

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \left(\frac{dz}{ds}\right)_{s=0} \int_{-\infty}^{+\infty} e^{-\Omega s^2} ds$$

(This gives the leading term of the asymptotic expansion.)

Hence

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \left(\frac{dz}{ds}\right)_{s=0} \sqrt{\frac{\pi}{\Omega}}$$

To evaluate the derivative:

$$-s^{2} = u(z) - u(z_{0})$$

$$= g(z) - g(z_{0}) \quad \text{(Recall: } v \text{ is constant along SDP.)}$$

$$-2s\left(\frac{ds}{dz}\right) = g'(z)$$
 At the saddle point this gives $0 = 0$.

Take one more derivative:

$$-2\left(\frac{ds}{dz}\right)\left(\frac{ds}{dz}\right) - 2s\frac{d^2s}{dz^2} = g''(z)$$

At
$$s = 0$$
:
$$\left(\frac{ds}{dz}\right)_{z_0} = \left(\frac{g''(z_0)}{-2}\right)^{1/2}$$

$$\left(\frac{dz}{ds}\right)_{s=0} = \left(\frac{-2}{g''(z_0)}\right)^{1/2}$$

Hence, we have

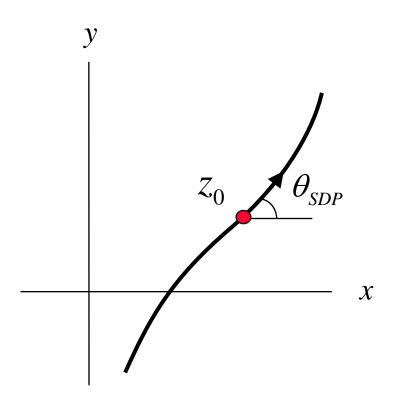
$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \left(\frac{-2}{g''(z_0)}\right)^{1/2} \sqrt{\frac{\pi}{\Omega}}$$

Note: There is an <u>ambiguity</u> in sign for the square root:

$$\left(\frac{dz}{ds}\right)_{s=0} = \left(\frac{-2}{g''(z_0)}\right)^{1/2}$$

To avoid this ambiguity, define

$$\arg\left(\frac{dz}{ds}\right)_{s=0} = \theta_{SDP}$$



The derivative term is therefore

$$\left(\frac{dz}{ds}\right)_{s=0} = \sqrt{\frac{2}{\left|g''(z_0)\right|}} e^{j\theta_{SDP}}$$

Hence

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \sqrt{\frac{\pi}{\Omega}} \sqrt{\frac{2}{|g''(z_0)|}} e^{j\theta_{SDP}}$$

To find θ_{SDP} :

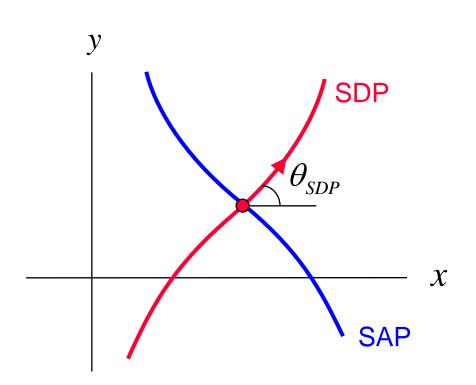
Denote:

$$g''(z_0) = R e^{j\alpha}$$
 $\alpha = \arg(g''(z_0))$
 $z - z_0 = r e^{j\theta_{SDP}}$

$$g(z) - g(z_0) \approx \frac{1}{2} g''(z_0) (z - z_0)^2 = \frac{1}{2} (Rr^2) e^{j(\alpha + 2\theta_{SDP})}$$

$$\alpha + 2\theta_{SDP} = \pm \pi$$

$$\theta_{SDP} = -\frac{\alpha}{2} \pm \frac{\pi}{2}$$



$$\theta_{SDP} = -\frac{\alpha}{2} \pm \frac{\pi}{2}$$

$$\alpha = \arg(g''(z_0))$$

Note: The direction of integration determines The sign.

The "user" must determine this.

Summary

$$I(\Omega) = \int_{C} f(z)e^{\Omega g(z)}dz$$

$$I(\Omega) \sim f(z_0) e^{\Omega g(z_0)} \sqrt{\frac{\pi}{\Omega}} \sqrt{\frac{2}{|g''(z_0)|}} e^{j\theta_{SDP}}$$

$$\theta_{SDP} = -\frac{\alpha}{2} \pm \frac{\pi}{2}$$

$$\alpha = \arg(g''(z_0))$$

Example

$$J_{0}(\Omega) = \frac{1}{\pi} \int_{-\pi/2}^{\pi/2} \cos(\Omega \cos z) dz$$
$$= \frac{1}{\pi} \operatorname{Re} I(\Omega)$$

where

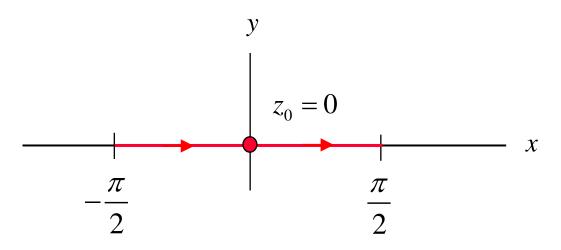
$$I(\Omega) = \int_{-\pi/2}^{\pi/2} e^{\Omega(j\cos z)} dz$$

Hence, we identify:

$$f(z) = 1$$
$$g(z) = j\cos z$$

$$g'(z_0) = -j\sin z_0 = 0$$

$$z_0 = 0 \pm n\pi$$



$$g(z) = j\cos z$$

$$g''(z_0) = -j\cos z_0 = -j$$

$$\alpha = \arg g''(z_0) = -\frac{\pi}{2}$$

$$\theta_{SDP} = -\frac{\alpha}{2} \pm \frac{\pi}{2} = \frac{\pi}{4} \pm \frac{\pi}{2}$$

$$\theta_{SDP} = -\frac{\pi}{4} \text{ or } \frac{3\pi}{4}$$

Identify the SDP and SAP:

$$g(z) = j\cos(x + jy)$$

= $j[\cos x \cosh y - j\sin x \sinh y]$

$$u(x, y) = \sin x \sinh y$$

 $v(x, y) = \cos x \cosh y$

SDP and SAP:
$$v(z) = v(z_0) = \text{constant}$$

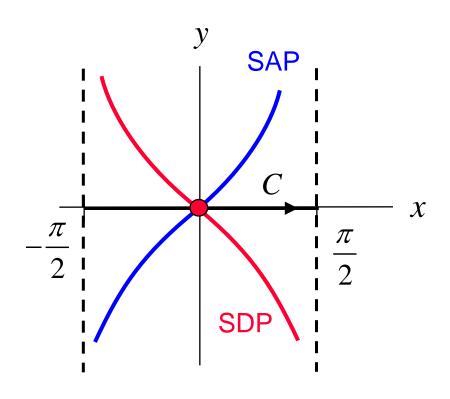
$$\cos x \cosh y = \text{constant} = \cos(0)\cosh(0) = 1$$

SDP and SAP:

 $\cos x \cosh y = 1$

Examination of the *u* function reveals which of the two paths is the SDP.

$$u(x, y) = \sin x \sinh y$$



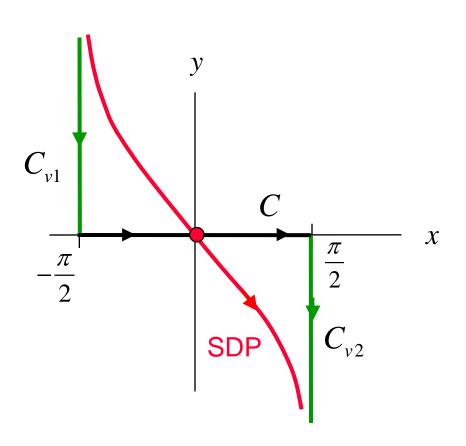
Vertical paths are added so that the path now has limits at infinity.

$$SDP = C + C_{v1} + C_{v2}$$

It is now clear which choice is correct for the departure angle:

$$\theta_{SDP} = -\frac{\pi}{4}$$

$$I(\Omega) = I_{SDP} - I_{v1} - I_{v2}$$



$$I_{SDP} \sim f(z_0) e^{\Omega g(z_0)} \sqrt{\frac{\pi}{\Omega}} \sqrt{\frac{2}{\left|g''(z_0)\right|}} e^{j\theta_{SDP}}$$

This is the answer for $I(\Omega)$ if we ignore the contributions from the vertical paths.

Hence,
$$I(\Omega) \sim (1) e^{\Omega j \cos(0)} \sqrt{\frac{\pi}{\Omega}} \sqrt{\frac{2}{|-j|}} e^{-j\frac{\pi}{4}}$$

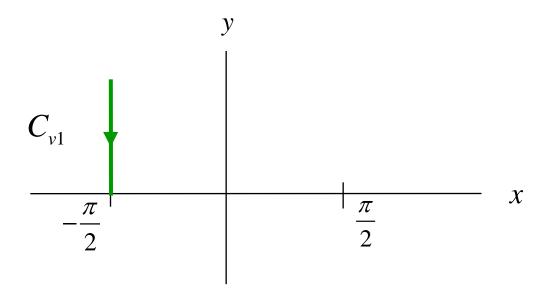
or
$$I(\Omega) \sim \sqrt{\frac{2\pi}{\Omega}} e^{j\left(\Omega - \frac{\pi}{4}\right)}$$

So
$$J_0(\Omega) \sim \frac{1}{\pi} \operatorname{Re} \left\{ \sqrt{\frac{2\pi}{\Omega}} e^{j\left(\Omega - \frac{\pi}{4}\right)} \right\}$$

Hence

$$J_0(\Omega) \sim \sqrt{\frac{2}{\pi \Omega}} \cos \left(\Omega - \frac{\pi}{4}\right)$$

Examine the path C_{v1} (the path C_{v2} is similar).



$$I_{v1} = \int_{-\pi/2 + j\infty}^{-\pi/2} e^{\Omega(j\cos z)} dz$$

Let
$$z = -\frac{\pi}{2} + jy$$

$$I_{v1} = j \int_{-\infty}^{0} e^{j\Omega \cos\left(-\frac{\pi}{2} + jy\right)} dy = j \int_{-\infty}^{0} e^{-\Omega \sinh y} dy = -j \int_{0}^{\infty} e^{-\Omega \sinh y} dy$$

$$\operatorname{since} \cos\left(-\frac{\pi}{2} + jy\right) = \sin(jy) = j \sinh y$$

Use integration by parts (we can also use Watson's Lemma):

$$I_{v1} = -j \int_{0}^{\infty} \left(\frac{-1}{\Omega \cosh y} \right) \frac{d}{dy} \left(e^{-\Omega \sinh y} \right) dy$$

$$\sim -j \left(\frac{-1}{\Omega \cosh y} \right) \left(e^{-\Omega \sinh y} \right) \Big|_{0}^{\infty}$$

$$I_{v1} \sim -j \left(\frac{1}{\Omega} \right)$$

Hence,

$$I_{v1} = O\left(\frac{1}{\Omega}\right)$$

Note: I_{v1} is negligible compared to the saddle-point contribution as $\Omega \to \infty$.

However, if we want an asymptotic expansion that is accurate to order $1/\Omega$, then the vertical paths <u>must</u> be considered.

Alternative evaluation of I_{v1} using Watson's Lemma (alternative form):

$$\begin{split} I_{v1} &= -j\int\limits_{0}^{\infty} e^{-\Omega \sinh y} dy \\ I_{v1} &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \left(\frac{dy}{ds}\right) ds \\ &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \frac{1}{\cosh(y)} ds \\ &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \frac{1}{\sqrt{1+\sinh^2(y)}} ds \\ &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \frac{1}{\sqrt{1+\sinh^2(y)}} ds \\ &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \frac{1}{\sqrt{1+sinh^2(y)}} ds \\ &= -j\int\limits_{0}^{\infty} e^{-\Omega s} \frac{1}{\sqrt{1+s^2}} ds \\ &= -j\int\limits_{$$

$$I_{v1} \sim -j \int_{0}^{\infty} e^{-\Omega s} (1) ds - j \int_{0}^{\infty} e^{-\Omega s} \left(-\frac{s^2}{2} \right) ds + \dots$$

SO

$$I_{v1} \sim -j \int_{0}^{\infty} e^{-\Omega s} ds + \frac{j}{2} \int_{0}^{\infty} e^{-\Omega s} s^{2} ds + \dots$$

Recall:

$$\int_0^\infty s^{\alpha_n} e^{-\Omega s} ds = \frac{1}{\Omega^{\alpha_n+1}} \Gamma(\alpha_n+1)$$

$$\Gamma(x) = (x-1)! \equiv \int_0^\infty t^{x-1} e^{-t} dt$$

Hence,

$$I_{v1} \sim -j\left(\frac{1}{\Omega}\right) + \frac{j}{2}\left(\frac{\Gamma(3)}{\Omega^3}\right) + \dots$$

Complete Asymptotic Expansion

By using Watson's lemma, we can obtain the <u>complete</u> asymptotic expansion of the integral in the steepest-descent method, exactly as we did in Laplace's method.

$$I(\Omega) = \int_{a}^{b} f(z) e^{\Omega g(z)} dz$$

Define:

$$s^{2} \equiv \left[g\left(z_{0}\right) - g\left(z\right) \right]_{SDP} \qquad h(s) \equiv f\left(z(s)\right) \left(\frac{dz(s)}{ds}\right)$$

Assume:
$$h(s) \sim \sum_{n=0,1,2,...} a_n s^n$$
 as $s \to 0$

Then we have

$$I(\Omega) \sim e^{\Omega g(z_0)} \sum_{n=0,2,4...} \frac{a_n}{\Omega^{\frac{(n+1)}{2}}} \Gamma\left(\frac{n+1}{2}\right)$$